

Article

Experiments on Cavitation Control around a Cylinder Using Biomimetic Riblets

Ebrahim Kadivar^{1,2} , Mazyar Dawoodian^{1,*} , Yuxing Lin¹ and Ould el Mactar¹ ¹ Institute of Ship Technology and Ocean Engineering, University of Duisburg-Essen, 47057 Duisburg, Germany; ebrahim.kadivar@uni-due.de (E.K.)² Faculty of Mechanical Engineering, Tarbiat Modares University, Tehran 1411713116, Iran

* Correspondence: mazyar.dawoodian@uni-due.de

Abstract: Experimental investigations were conducted to uncover the impact of cavitation control—through the use of biomimetic riblets on cavitating flows around a circular cylinder. First, the dynamics of cavitation in the flow behind a finite cylinder (without riblets) was unveiled by visualizing the cavitation clouds and measuring the lift force fluctuations acting on the cylinder. Second, in a significant step forward, a comprehensive explanation was provided for the cavitation control methods using two bio-inspired riblet morphologies positioned in different orientations and locations on the cylinder. For the first time, the impacts of these tiny formations on the flow dynamics and the associated cavitation process were scrutinized. This showed that scalloped riblets, with their curved design, induced secondary vortices near their tips and distorted primary streamwise vortices, and that high velocity gradients near the jagged pattern peaks of sawtooth riblets delayed flow separation, which affected cavitation.

Keywords: hydrodynamic cavitation; cavitation control; biomimetic riblets; cylinder



Citation: Kadivar, E.; Dawoodian, M.; Lin, Y.; el Mactar, O. Experiments on Cavitation Control around a Cylinder Using Biomimetic Riblets. *J. Mar. Sci. Eng.* **2024**, *12*, 293. <https://doi.org/10.3390/jmse12020293>

Academic Editor: Stefano Gaggero

Received: 9 January 2024

Revised: 27 January 2024

Accepted: 5 February 2024

Published: 6 February 2024



Copyright: © 2024 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<https://creativecommons.org/licenses/by/4.0/>).

1. Introduction

Cavitation appears as an undesirable phenomenon in diverse industrial applications. To elaborate, cavitation leads to erosion on the surface of propellers/rudders of ships, or it generates unwelcome vibration/noise within the components of hydraulic systems (see, e.g., Reisman et al. [1], Dular et al. [2], Haosheng [3], Patella et al. [4], Kadivar et al. [5], and Lin et al. [6]). Franc and Micheal [7] noted that, at a constant temperature, cavitation may occur around immersed objects when the local pressure falls below the vapor pressure of the liquid. Additional factors, including turbulent flow characteristics, surface roughness [8,9], temperature and viscosity of the liquid [10], and density and radius of the nuclei [11], may affect cavitation or contribute to its initiation on immersed bodies. Yuan et al. [12,13] compared the difference in ice breaking between a non-cavitating and a cavitating water jet, and they found that the collapse of cavitation produces additional pressure loads on the wall, which performs better in the later continuous damage process. Different types of cavitation commonly encountered on ship rudders and on the components of hydraulic systems manifest as sheet-like cavitation, partial cavitation, or cloud cavitation. Among these, cloud cavitation is the most hazardous because it tends to induce severe erosion damage, thereby provoking surface vibrations on the immersed bodies, as highlighted by Kadivar [14].

Prior investigations dealt with diverse cavitation dynamics that occur on hydrofoils. For instance, Li et al. [15] examined the partial cavitation dynamics in various regimes of cavitations on a hydrofoil surface. Callenaere et al. [16] investigated the partial cavitation occurring on a hydrofoil, along with the formation of a re-entrant jet that may arise within this specific cavitation regime. Their findings revealed that a re-entrant jet generates and progresses towards the leading edge of the hydrofoil in a regime of partial cavitation. This observation led them to conclude that this phenomenon is the predominant mechanism

responsible for the shedding of cavities from the surface of the hydrofoil. Meanwhile, Stutz and Reboud [17] conducted experimental analyses to dissect the sheet cavitation formations on a geometry such as Venturi. They scrutinized the behavior of attached cavitation, including the length patterns of cavities and the phenomenon of reversed flow. Leroux et al. [18,19] showed that two primary mechanisms contribute to the shedding of cavities on a hydrofoil surface in the partial cavitation regime. These mechanisms include the formation of a re-entrant jet and the formation of a shockwave subsequently causing the cavity's collapse. The shockwave alters the cavitation dynamics. Over recent decades, numerous investigations focused on comprehending the dynamics of cloud cavitation, specifically aiming to elucidate the mechanisms of cavity shedding and cavitation fluctuations. For instance, Ganesh et al. [20] explored the intricacies of unsteady cavitation emerging on an apex geometry through the use of a X-ray densitometer method. Their observations showed that a distinct shockwave contributes to the oscillation of unsteady cavitation at the aft part.

Ghahramani [21] showed that the small-scale cavities not only are important at the inception and collapse steps, but also influence the development of large-scale structures. Furthermore, they explain the effects of flow vorticity, pressure variation rate, bubble inertia and surface tension effects on the cavity dynamics. Kadivar et al. [22,23] investigated the unsteady cavitation dynamics and surge cavitation occurring over both a hydrofoil and a semicircular flat plate. They concluded that the evolution of the sheet cavity to a cloud cavity diverges from the cavity structure dynamics within the cloud cavitating regime. Their findings pointed out that within the surge cavitation regime, the cavity fluctuations' dynamics are driven by pressure waves generated within the cavity due to the collapsing of the cavity. In other words, the mechanism of the cavity shedding was not directly generated by a re-entrant jet in the cavitation surge regime. Long [24] investigates the cavitating flow around a Delft Twist-11 hydrofoil using the Zwart–Gerber–Belamri cavitation model combined with an LES turbulence model. They studied the behaviors of the re-entrant and side-entrant jets from both Euler and Lagrangian viewpoints, and they utilized the three-dimensional LCSs obtained with the three-dimensional Lagrangian technology to investigate the influence of U-type structures on local flow patterns. Pelz et al. [25] examined the influence of a re-entrant jet on unsteady cloud cavitation. Their research also presented insights into how the nucleation rate affects the cavity length and how the Reynolds number influences the formation of the re-entrant jet on the hydrofoil.

Wang et al. [26,27] present a multi-scale Euler–Lagrange method to provide deep insight into the unsteady instability mechanism and internal structures of cloud cavitation around the Delft Twist-11 hydrofoil. With their method, they analyzed the cloud cavitation evolution mechanism and the internal cavity structure, and they explained the bubble size distribution spectra in the unsteady cavitating flow. Circular cylinders or similar geometries are commonly employed in marine contexts. Extensive prior studies have addressed cavitation around circular cylinders and the mechanism of cavitation shedding. For instance, Matsudaira et al. [28] examined the flow characteristics within the vortex cavitation. They discovered that the vortices' shedding frequency aligns closely with the pressure pulsations caused by the collapse of the cavity. Saito and Sato [29] examined the formation of Karman vortex-like cavity structures behind a cylinder. They analyzed the cavitation pulsations and the erosion induced by the cavity collapse on a plate made of aluminum mounted to the test section's wall. Their hypothesis suggested that the erosion on the plate is caused by strong impulsive forces resulting from collapsing cavities.

Kumar et al. [30] studied experimentally the cavitation behind a cylinder at a Reynolds number of around 64,000. They revealed that reducing the cavitation number leads to a decrease in the frequency of cavity shedding. Geike [31] calculated macroscopic forces between moving rough surfaces, and the tensile stresses that affect the wear in mixed lubrication applications, and developed a model for the numerical simulation of cavitation in mixed lubrication contacts that reproduces tensile stresses on short time scales. Moreover, they identified that lower cavitation numbers correlate with higher pressure

pulsations. Gnanaskandan and Mahesh [32], by numerically analyzing the pressure waves induced by the cavitation collapses, illustrated the significant impact of cavitation on the pressure distribution, and the resultant forces. Dobroselsky [33] performed an experimental study of cavitating flows around a cylinder at critical Reynolds numbers. He found the quasi-periodic nature of the cavity shedding mechanism from the surface of the cylinder. Additionally, he indicated the dependence of cavitation onset on the Reynolds number.

Brandao et al. [34] performed a numerical simulation of the cavitation behind a cylinder at various cavitating regimes and low Reynolds numbers from the laminar to turbulent flows. Their findings revealed a significant alteration in cavitation within the Karman vortices caused by the initiation of the propagation of condensation front. Gu et al. [35] employed the Large Eddy Simulation (LES) technique to explore cavitation dynamics behind a cylinder at a Reynolds number of around 9500. Their findings highlighted that vortex cavitation plays an important role in influencing the shedding frequency of the cylinder's wake. Sadri and Kadivar [36] investigated the behavior of unsteady cavitation and the noise generated by the cavity collapses around one and two cylinders. They specifically analyzed the cavitation in scenarios where two cylinders were positioned side by side with varying gaps between them. Their study indicated that, through gap reduction, the wake patterns formed behind these cylinders coalesce, giving rise to a singular vortex street. Zhang [37] established a novel theory for oscillating bubble dynamics, which can simultaneously consider complex physical factors such as boundaries, bubble interaction, ambient flow field, gravity, bubble migration, fluid compressibility, viscosity, and surface tension. Their theory achieves the unification of classical bubble theories. They showed that their theory has higher accuracy and applicability compared to the classical ones, and may provide new references for future explorations in bubble dynamics, cavitation, and other multiphase flow-related problems.

Regarding cavitation control, extensive studies were conducted to explore various passive methods aimed at mitigating cavitation effects around immersible bodies. Custodio et al. [38] examined the impact of protuberances on cavitation. Their findings indicated that medium to large protuberance amplitudes hinder cavitation development at the aft part of the protuberance troughs. Hao et al. [39] analyzed the dynamics of cavitation around hydrofoils with different surfaces and they emphasized that the presence of a roughness across the hydrofoil's surface may affect the cloud cavitation. They deduced that hydrofoils with a rough surface pattern experience heightened intensity of cloud cavitation. Zhao and Wang [40] investigated unsteady cloud cavitation control by implementing a bio-inspired structure situated on the suction side of the hydrofoil, showing that the bio-inspired structure possesses the capability to control the levels of turbulent kinetic energy. Furthermore, they demonstrated that hydrodynamic efficiency improves due to the presence of the bio-inspired structure. Zhang et al. [41] conducted both experimental and numerical studies focused on controlling the cavitation by incorporating an obstacle on the hydrofoil. They showed that the obstacle restricts the cavity fluctuations.

Che et al. [42,43] investigated the control of partial cavitation on a hydrofoil through the use of micro vortex generators. Their findings revealed that these micro vortex generators effectively shifted the location of cavitation inception closer to the point of flow separation within the hydrofoil's unsteady cavitating regime. Kadivar et al. [44,45] showed that utilizing bubble generators effectively mitigates cavitation. They also conducted cavitation control on a hydrofoil surface, employing passive methods including cylindrical and wedge-type vortex generators. Their study demonstrated effective suppression of cloud cavitation on the hydrofoil with the vortex generators, leading to a substantial reduction in high-pressure pulsations resulting from the cloud cavitation collapse. Simanto et al. [46] showed that leading-edge protuberances mitigate the cavitation, and these protuberances reduce the noise, particularly at higher Reynolds numbers when the cavitation is the primary source of noise. Yu et al. [47] experimentally and numerically studied controlling cavitation using a porous pattern on a cylinder, demonstrating that a porous layer mitigated the generated vortices and the inception of cavities.

To the authors' knowledge, the current work presents a unique experimental analysis that goes deeper than previous analyses to effectively reveal the detailed mechanism of cavitation control using riblets. The objective was to significantly enhance our understanding of cavitation dynamics and control mechanisms. The paper was organized as follows. Section 2 describes the experimental setup, the experimental conditions, and the test cases considered. Section 3 portrays the main results, including the dynamics of cavitation in the flow region behind a cylinder, the morphological analysis of riblets, and the effects of flow conditions on cavitation control. Finally, Section 4 summarizes the conclusions.

2. Experimental Setup

The investigations were performed in the cavitation tunnel K23 at the Institute of Ship Technology of the Duisburg-Essen University. This tunnel functioned as a closed cavitation tunnel, comprising an electric motor-driven impeller, a vacuum pump, and an array of pressure transducers and temperature sensors. The tunnel reached inlet velocities of up to 9.0 m/s, with fluctuations of typically less than 0.1 m/s. The tunnel's test section (1.1 m in length) encompasses a square inlet section of 0.3 by 0.3 m². This test section, consisting of four Plexiglas walls, offered optimal observation from various perspectives. A high-speed camera, positioned at one side of the test section, monitored the flow. During the experiments, cylinders were vertically installed on the test section wall. Figure 1 illustrates the experimental setup, and the top view of two different configurations of riblet placement, with an arrow pointing in the flow direction.

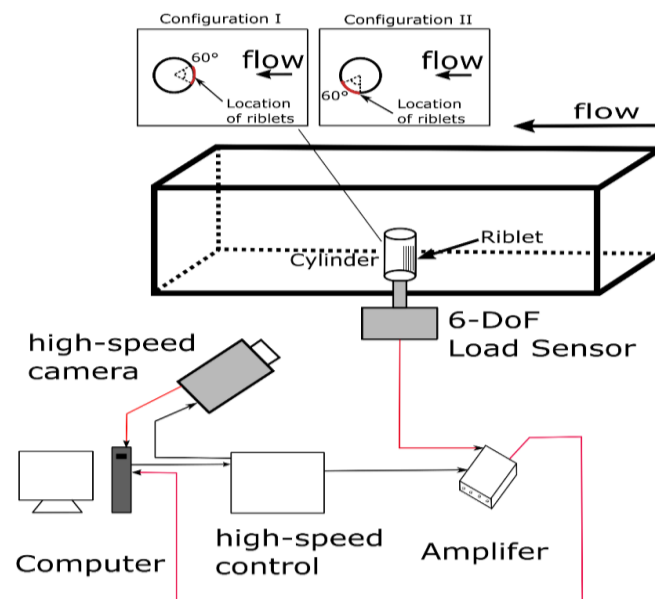


Figure 1. Schematic plan view of the test section with instrument. The top view presents two different configurations of riblet placement, with an arrow pointing in the flow direction.

A high-speed camera was used to capture the cavitation around a cylinder with and without riblet structures. A six DOF freedom load sensor, located on one side of the tunnel's test section between the tunnel wall and the cylinder, measured the lift and drag forces acting on the circular cylinder. A pressure transducer, mounted at the outlet of the tunnel, measured the pressure fluctuations behind the cylinder and in the cylinder's wake. The experiments were conducted at different Reynolds numbers: $Re = 1.0 \times 10^5$, 1.25×10^5 , and 1.5×10^5 ; and at three cavitation numbers of $\sigma = 1.0$, 1.5 , and 2.0 . The cavitation number σ [2] and the Reynolds number Re were here defined as follows:

$$\sigma = \frac{p_{ref} - p_v}{1/2 \rho_{ref} V_{ref}^2} \quad (1)$$

$$\mathfrak{R} = \frac{V_{ref} l_{ref}}{\nu} \quad (2)$$

where p_{ref} and p_v are the reference static pressures at the middle of the inlet section and the saturation vapor pressure of the operating liquid, respectively; V_{ref} is the inlet flow velocity; and ρ_{ref} and ν are the reference density and kinematic viscosity of the water, respectively. The experimental conditions were constantly adjusted during the tests to assess the cavitation dynamics between comparative cases. To reduce the effect of varying air content, the air inside the cavitation tunnel was kept within the range of 1.2 mg/L and the water temperature was about 16 °C with a tolerance of 0.1 °C. Consequently, the uncertainty of the vapor pressure was almost 24 Pa; the uncertainty of the inlet pressure ranged between 100 and 200 Pa; the uncertainty of the cavitation number ranged between 0.05 and 0.10; and the uncertainty of the Reynolds number was between 1 and 2%. The high-speed imaging was captured at a frequency of 1000 Hz with an exposure time of 30 μ s. The resolution of the captured images was 1632 $\times$ 1200 pixel counts. A sampling rate of about 5 Hz was specified for the acquisition system that measures the force. Employing a measuring amplifier MX1601B, the six DoF force sensor relied on an HBM data acquisition system to capture the force fluctuations on the cylinder. A LaVision high-speed control system enabled the acquisition and post-processing of the captured images. A 5V trigger signal from this control system to the HBM measuring amplifiers synchronized the captured images with the data acquisition system.

Figure 2 shows the cylinders with horizontal/vertical scalloped/sawtooth riblets. The 13.0 cm long cylinders had a diameter of 2.5 cm. The riblets on the manufactured cylinders were 1.5 mm in height and the spacing between riblets measured 3.0 mm. The top width of the scalloped riblets was 0.5 mm, and the groove radius between these riblets was 1.0 mm. The section of the manufactured riblets on the cylinder was 8.0 cm, that is, the riblets extended over only one part of the cylinder's circumference and not over the rest of the cylinder.



Figure 2. (a) Schematic view of the cylinder with vertical and horizontal sawtooth riblets and photos of the manufactured sawtooth riblets on the cylinders. (b) Schematic view of the cylinder with vertical and horizontal scalloped riblets and photos of the manufactured scalloped riblets on the cylinders.

3. Results and Discussion

3.1. Cavitation Dynamics

Figure 3 presents the cavity dynamics behind the cylinders caused by the flow at $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$. Figure 3a illustrates the cylinder without riblets. Figure 3b,c depict the cylinder with horizontal and vertical sawtooth riblets, respectively. Similarly, Figure 3d,e represent the cylinder equipped with horizontal and vertical scalloped riblets, respectively. As can be seen, the flow interacted with the cylinder, and miniature riblets induced the creation of small vortex structures nearby, where they significantly affected any instabilities caused by the absence of the riblets. In the scenario in which a smooth cylinder (without riblets) of finite length and with a free-end (tip) was exposed to an incident flow, a pair of counter-rotating trailing vortices emerged near the cylinder's tip, as shown by the purple-colored vectors in Figure 3a. The formation of these vortices arose due to a phenomenon termed upwash; this is caused by a weak upward-directed local velocity field at the cylinder's tip [48]. As the flow proceeded over the tip, a downwash occurred, whereby the flow descended into the low-pressure region immediately behind the cylinder. The combination of an upwash flow on the sides and the downwash at the center triggered the formation of two counter-rotating trailing vortices [49]. Additionally, a segment of the incident flow close to the base moved downward as it approached the cylinder [50], and it subsequently circulated upstream of the junction between the cylinder and the wall. In this region, the horseshoe vortex took on its shape, wrapped around the cylinder, as shown by the yellow vectors in Figure 3a, and then progressed downstream [51]. The combined presence of upwash and downwash vortices led to shedding patterns of varying frequency and amplitude and generated unstable Karman vortices, as shown by the white vectors in Figure 3a, and this intricate, dancing vortex trio added another layer of complexity to the formation of cavitation.

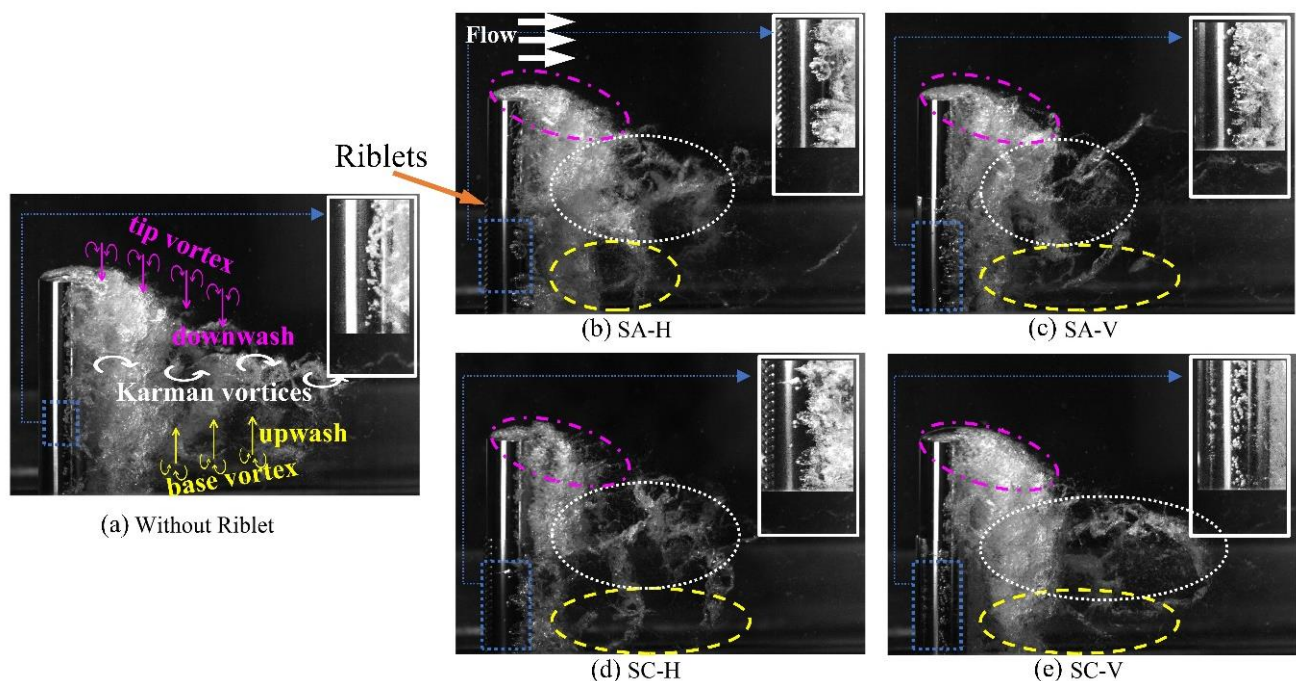


Figure 3. At $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$, cavitation patterns behind the circular cylinder without riblets (a), with horizontal sawtooth riblets SA-H (b), with vertical sawtooth riblets SA-V (c), with horizontal scalloped riblets SC-H (d), and with vertical scalloped riblets SC-V (e). As the flow interacted with the cylinder, miniature riblets induced the creation of small vortex structures nearby, and these small vortices significantly impacted the instabilities that typically formed on the plain cylinder's surface.

Cavitation on the cylinder occurred akin to hydrofoil cavitation, particularly in the proximity of the laminar separation bubble (LSB) and close to the flow reattachment point [23]. It started with the generation of linear instability-driven Tollmien–Schlichting (T–S) waves, i.e., constant-frequency traveling waves within the boundary layer. These T–S waves in particular contributed to laminar separation bubble formation. Subsequently, nonlinear three-dimensional (3D) instabilities arose from these linear waves that led to turbulence. Turbulent spots formed and expanded in both streamwise and spanwise directions. This T–S wave development within the stable laminar boundary layer crucially triggered the transition to turbulent flow. Pressure fluctuations, caused by flow instabilities on the cylinder surface, induced cavitation when the local water pressure dropped below the vapor pressure [52]. We focused our cavitation control strategy on managing flow instabilities, as well as pressure fluctuations, and mitigating LSB formation using miniature riblets.

3.2. Morphology of Riblets

Riblets come in a range of shapes, which are often inspired by biological forms, and they are commonly employed to manage flow separation and reduce drag [53–55]. The most common riblets include sawtooth and scalloped riblets. The usefulness of riblets is based on their capacity to generate vortices behind them, enabling control of the boundary layer. As shown in Figure 3b–e, the geometric configuration of riblets and the method of their implementation on the cylinder’s surface significantly influenced flow control and, thus, the inception of cavitation. When considering flow control, the choice of riblet type plays a significant role. Optimal selection involves selecting riblets capable of directing fluid with higher momentum to the boundary layer, where momentum is comparatively low. This strategic selection aids in enhanced control of the boundary layer; a critical factor for effective flow management.

The mechanism behind controlling cavitation achieved through the utilization of riblets has been elucidated in our prior published research [44,45]. As the flow interacts with the cylinder, the presence of miniature riblets gives rise to the formation of small vortex structures in their vicinity. These compact vortices play a pivotal role in influencing the instabilities present within the boundary layer, which would otherwise emerge on the cylinder’s surface in the absence of riblets. This mechanism operates through the generation of streamwise vortices by the riblets, facilitating the transfer of fluid with higher momentum from the upper layer into the boundary layer. A key outcome of implementing riblets is the enhancement of boundary layer velocities near the cylinder surface, leading to a reduction in velocity pulsations within the boundary layer. Thus, employing riblets enabled us to manage the large-scale cavity pulsations generated behind the cylinder. This was achieved by decreasing the instabilities in the upstream boundary layer on the cylinder surface.

Figure 4 presents flow patterns developed near the scalloped riblet’s tip. As can be seen, scalloped riblets altered the roll-up and breakdown of streamwise vortices in different ways, leading to decisive changes in behavior and interaction with the surrounding flow. Thus, their effects varied, based on their ability to generate vortices influencing the boundary layer. The presence of scalloped riblets resulted in the generation of secondary vortices [56] inside the valleys near the tip of the riblets (see Figure 4a). In certain instances, both the primary large-scale streamwise vortices and these secondary vortices, induced by the riblet tips, underwent distortion (as seen in Figure 4b) and acquired an elongated shape. This altered vortex morphology significantly affected the flow dynamics and cavitation inception. These alterations in vortex structure hold substantial implications for flow dynamics.

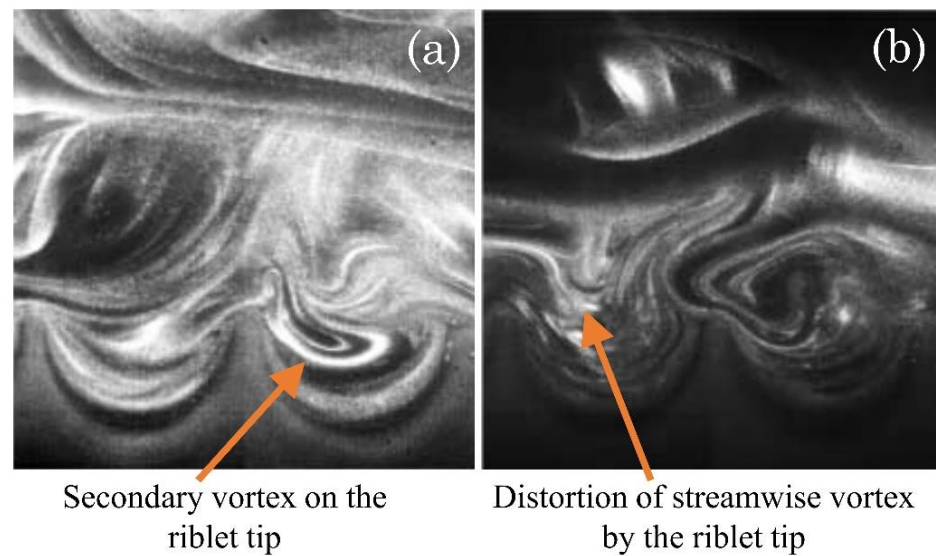


Figure 4. Scalloped riblets induced secondary vortices (a). In certain instances, both the primary large-scale streamwise vortices and these secondary vortices underwent distortion and became elongated (b). Reproduced with permission from the journal *Experiments in Fluids* [56]; published by Springer Nature (2024).

In the case of sawtooth riblets, Goldstein et al. [57] observed a distinctive flow behavior. Close to the peaks of sawtooth riblets, the velocity gradient increased notably, suggesting a region of high velocity gradients. Conversely, within the valleys between the riblets, the velocity gradient is notably reduced. Recall that, occasionally, vorticities do extend into these valleys, but they are usually concentrated near the tips of the riblets. However, as seen in Figure 3b,c, while the jagged pattern of sawtooth riblets actively controlled the near-wall flow by delaying flow separation, the curved pattern of scalloped riblets depicted in Figure 3d,e delayed the transition to turbulence.

Horizontal riblets modified the near-wall flow, predominantly affecting turbulence levels, while vertical riblets smoothed the flow. Figure 3b,d demonstrate that horizontal riblets promote a stable and predictable flow. In contrast, Figure 3c,e illustrate that vertical riblets alter the intensity and behavior of turbulence structures. Thus, riblets with different patterns, i.e., scalloped or sawtooth mounted horizontally or vertically, offered unique ways to alter flow dynamics [58] as well as the associated turbulence and separation characteristics of the Karman vortices. Consequently, it is crucial to select riblets capable of redirecting the fluid with higher momentum into the boundary layer, even when the flow's momentum is relatively low.

Figures 5 and 6 present five sequential images at five time intervals (periodic) depicting the cavitation process behind the cylinder with horizontal scalloped (SC-H) riblets and vertical scalloped (SC-V) riblets, both caused by the flow at $\text{Re} = 1.25 \times 10^5$ and $\sigma = 1.0$. These specific images were chosen from half of a standard period T (cyclic) that characterized the dynamics of cavity shedding. The red areas indicate regions where the riblets exerted their most significant influence on cavitation. Notably, the pressure increasing behind the riblets led to a reduction in the cavity volume behind the plain cylinder. As expected, in the case of SC-H riblets, an effective vortex interaction with the boundary layer is seen (Figure 5). This occurred because flow moved through the tips and valley curves of the horizontal (SC-H) riblets, generating strong longitudinal vortices that affected the boundary layer. In comparison, the flow over the vertical scalloped riblets (SC-V) primarily engaged with the riblets' sharp edges (Figure 6), promoting the merging of vortices further downstream.

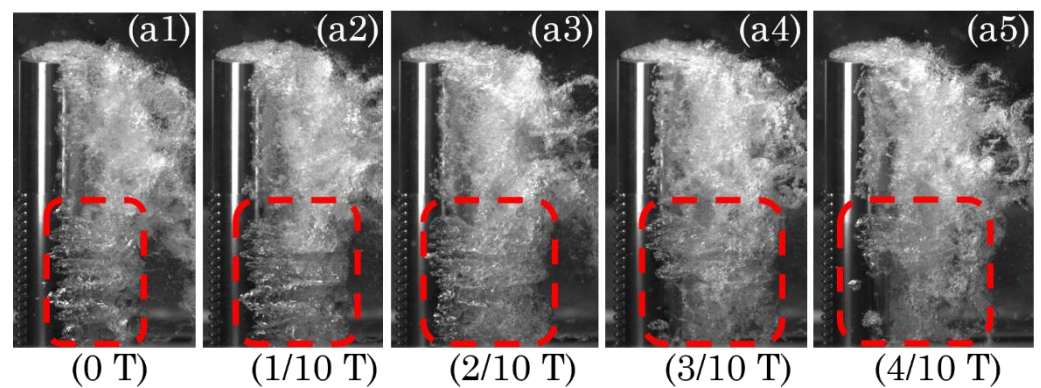


Figure 5. At $\mathcal{R} = 1.25 \times 10^5$ and $\sigma = 1.0$, five sequential images captured at regular intervals, showing the progression of cavitation behind the cylinder with SC-H riblets. These specific images were chosen from half of a standard cyclic period T that characterized shedding of cavitation. Red surrounded areas indicate regions where the riblets exerted their most significant influence on cavitation. The pressure increasing behind the SC-H riblets reduced the cavity volume when compared to this volume behind the plain cylinder. The images (a1), (a2), (a3), (a4) and (a5) show the cavitation dynamics at time steps $0 T$, $1/10 T$, $2/10 T$, $3/10 T$, $4/10 T$, respectively.

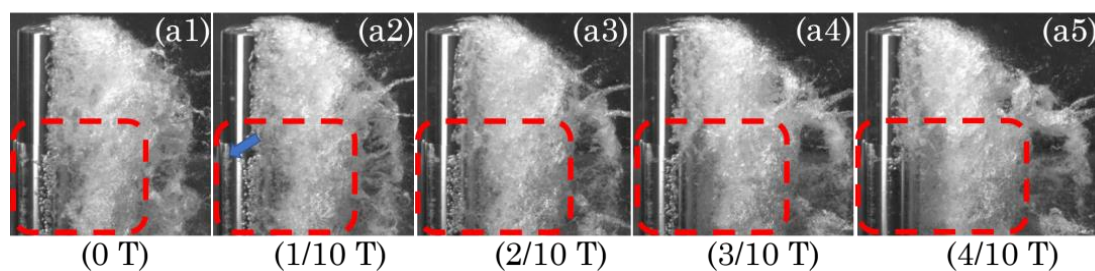


Figure 6. At $\mathcal{R} = 1.25 \times 10^5$ and $\sigma = 1.0$, five sequential images captured at regular intervals, showing the progression of cavitation behind the cylinder with SC-V riblets. These images were chosen from half of a standard cyclic period T that characterized shedding of cavitation. Red surrounded areas indicate regions where the riblets exerted their most significant influence on cavitation. The flow over SC-V riblets primarily engaged with the sharp edges of the riblets, promoting the merging of vortices further downstream. The images (a1), (a2), (a3), (a4) and (a5) show the cavitation dynamics at time steps $0 T$, $1/10 T$, $2/10 T$, $3/10 T$, $4/10 T$, respectively.

For the cylinder with SC-V riblets, depicted in Figure 3e, the flow around these riblets mostly interacted with the sharpness of the riblets. Instead of longitudinal vortices, this resulted in a vortex sheet forming behind the riblets. SC-H riblets facilitate better interaction as flow moves through valleys at both sides of the riblet instead of merely passing over them, as with SC-V riblets. This increased interaction along the top and bottom surfaces of each riblet led to the merging of vortices downstream, and created stronger vortices. Consequently, the horizontal (SC-H) riblets effectively directed higher-momentum fluid into the boundary layer due to their enhanced interaction with the flow.

We also conducted a more comprehensive analysis of how riblet shape and orientation impacted cavitation dynamics in flows at constant Reynolds and cavitation numbers. Riblets modified large-scale cavity structures into smaller ones. Consequently, the cylinder with riblets anticipated lower pressure fluctuations, resulting from detachment and also collapse of these smaller cavities. In essence, riblets mitigated lift force fluctuations and cavitation-induced vibrations on the cylinder surface, resulting in reduced amplitude and shedding frequency. This represents a significant advantage of riblet implementation. Note that the cavity shedding frequency was slightly higher for the cylinder with riblets because the new small-scale cavities generated by these riblets had a higher frequency of cavity shedding compared to the large-scale cavities.

Figure 7 presents comparative distributions of lift force F_y versus frequency f for the cylinder without and with riblets, caused by the flow at $R = 1.25 \times 10^5$ and $\sigma = 1.0$. Figure 7a plots these distributions for the cylinder with SC-H and SA-H riblets; Figure 7b does so for the cylinder with SC-H and SC-V riblets. The frequencies with highest amplitudes corresponded to the shedding frequency of large-scale cavity structures.

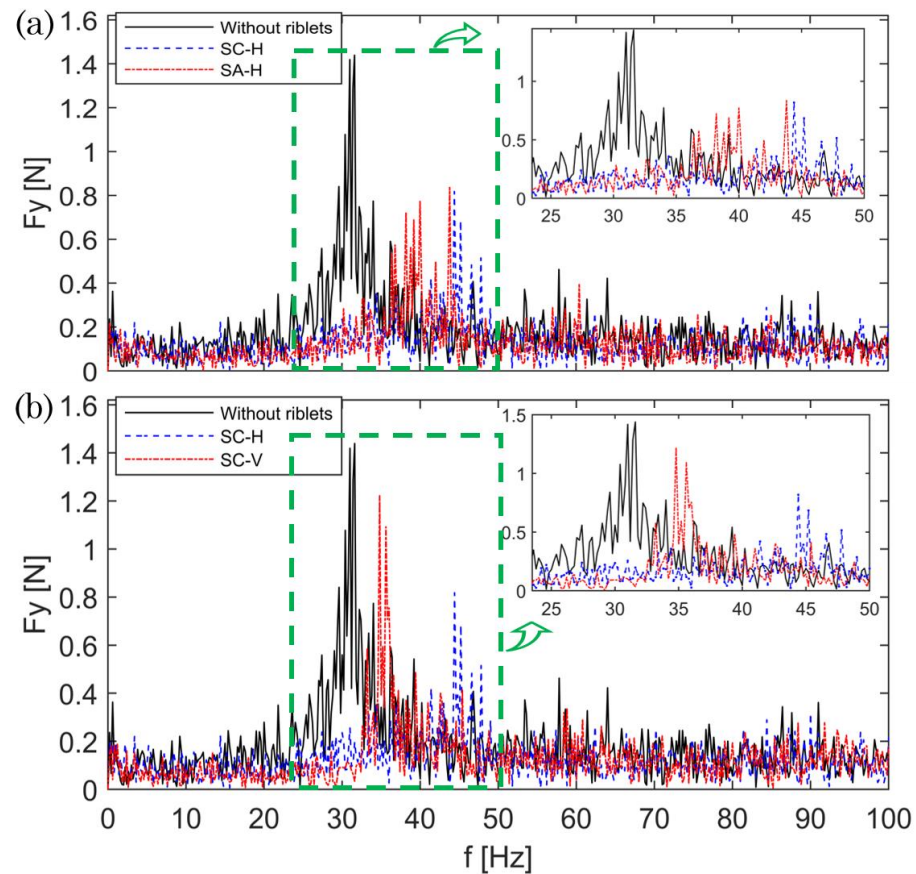


Figure 7. At $R = 1.25 \times 10^5$ and $\sigma = 1.0$ (a) distributions of the frequency based on the lift force measurement for the plain cylinder, SC-H, and SA-H cylinders; and (b) for the plain cylinder, and SC-H and SC-V cylinders. Frequencies with the highest amplitudes corresponded to the shedding frequency of large-scale cavity structures. For cylinders with riblets, lift force amplitudes were reduced by 41 and 43% for SA-H and SC-H cylinders, respectively; and by 43 and 30% for the SC-H and SC-V cylinders, respectively.

As can be seen in Figure 7a, when the riblets were mounted horizontally, the amplitude of the lift force was substantially reduced; specifically, by 41% (from 1.44 to 0.84 N) for SA-H, and by 43% (from 1.44 to 0.82 N) for SC-H. Furthermore, the shedding frequency of cavities increased for SA-H and SC-H cases compared to the plain cylinder; specifically, from 31.6 Hz to 43.8 Hz. In Figure 7b, for the SC-H cylinder, the lift force was also reduced by 43% (from 1.44 to 0.82 N), and by 30% (from 1.44 to 1.01 N) for SC-V. Here too, the shedding increased for the cylinder with scalloped riblets (SC-H and SC-V) compared to the smooth cylinder; specifically, from 31.6 to 44.4 Hz. Remarkably, horizontal scalloped riblets (SC-H) reduced the amplitude of lift fluctuations more effectively than vertical scalloped riblets (SC-V).

Figure 8 illustrates a comparison between the lift force amplitude and cavitation shedding frequency for plain cylinders and those with riblets mounted horizontally and vertically. The findings indicate that, in the case of scallop riblets, horizontal installation results in a smaller lift force range. Conversely, for sawtooth riblets, vertical implementation has a more pronounced impact on reducing the amplitude of lift forces. Nonetheless, the

shedding frequency is higher for cylinders with horizontal scalloped/sawtooth riblets compared to their vertical counterparts, with the cylinders without riblets having the lowest shedding frequency. One reason for this increase in shedding frequency is the reduced time required for cavities to detach from the surface of the cylinder due to the presence of riblets. Consequently, riblet structures can be inferred to reduce the amplitude of lift fluctuations while slightly increasing the shedding frequency.

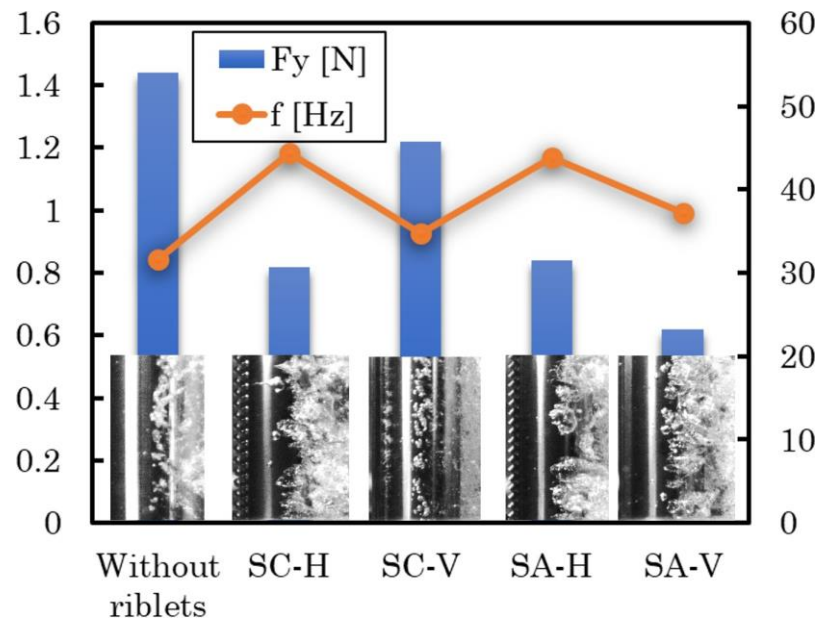


Figure 8. At $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$, lift force (blue columns, left scale) and vortex shedding frequency (coral dots, right scale) for the plain cylinder, SC-H, SC-V, SA-H, and SA-V cylinders. Horizontal scalloped riblets and vertical sawtooth riblets led to smaller lift force amplitudes. The shedding frequency of the cavity for SC-H and SA-H cylinders was higher compared to the frequency of the cavity shedding for the cylinders with vertical riblets, while the shedding frequency was lowest for the plain cylinder.

Although placing riblets on the upstream face of the cylinder mostly altered the attached flow and the flow's stability, rear riblets influenced the shedding of vortices in the wake region, thereby also mitigating cavitation. To explore the effects of placing riblets on the cylinder's rear side, the cylinder with SC-H riblets was rotated 150 deg clockwise. Figure 9 shows the riblets situated on the upstream/downstream face of the cylinder, caused by the flow at $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$. As can be seen, the riblets affected the attached flow and the flow's stability, whereas rear riblets impacted vortex shedding in the wake region, potentially mitigating cavitation. The red surrounded area in Figure 9a shows the extensive cavitation cloud generated behind the plain cylinder as this cloud was transformed into smaller and less dense clouds for the cylinder with horizontal scalloped riblets (SC-H) facing upstream. The green surrounded areas in Figure 9b,c represent these cavitation clouds for the cylinder with unrotated and rotated SC-H riblets. While upstream-faced SC-H riblets promoted a smoother flow over the cylinder's surface and delayed the transition to turbulence, the rotated SC-H riblets generated unstable and unpredictable flow patterns. These flow patterns reduced the shedding frequency of cloud cavitation.

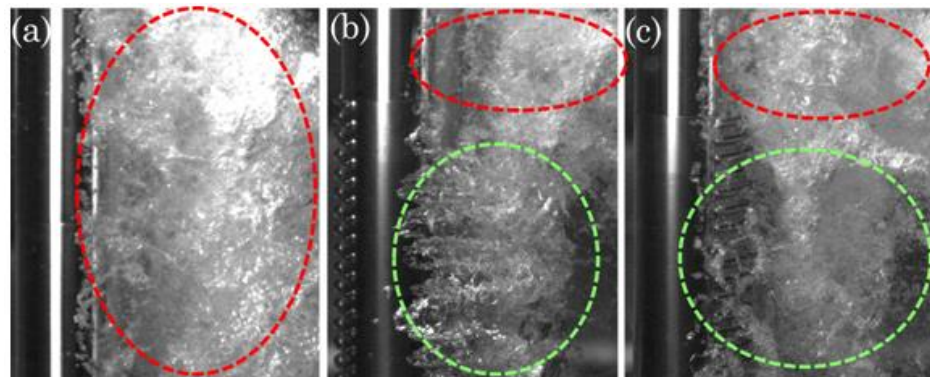


Figure 9. At $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$, placing riblets on the upstream face of the cylinder affected the attached flow and the flow's stability, whereas rear riblets impacted vortex shedding in the wake region, thereby mitigating cavitation. The extensive cavitation cloud observed behind the smooth cylinder (a), indicated by the red surrounded area, was transformed into smaller and less dense clouds for the cylinder with upstream-faced SC-H riblets (b) and for the cylinder with rotated SC-H riblets (c), as indicated by the green surrounded areas. Although upstream-faced SC-H riblets promoted a smoother flow over the cylinder's surface and delayed the transition to turbulence, rotated SC-H riblets generated unstable and unpredictable flow patterns.

Figure 10 presents five sequential images, captured at regular intervals, showing the progression of cavities around a cylinder with SC-H riblets rotated 150 deg clockwise, caused by the flow at $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$. While the upstream-oriented SC-H riblets resulted in a smoother flow and delayed the transition to turbulence, the rotated SC-H riblets induced unstable and unpredictable flow patterns. Although the rotated riblets had a lesser effect on modifying the wake behavior, they caused unstable and unpredictable flow patterns.

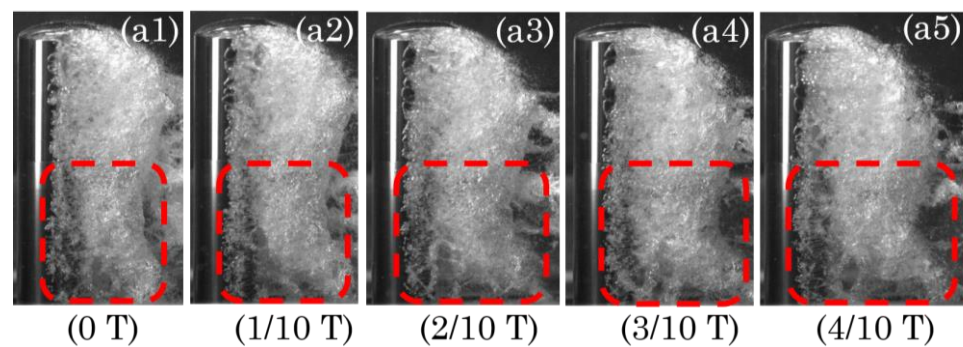


Figure 10. At $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$, five sequential images captured at regular intervals, showing the progression of cavitation dynamics around a cylinder with rotated (150 degree clockwise) SC-H riblets. These images were chosen from half of a standard cyclic period T that characterized cavitation shedding dynamics. Red surrounded areas indicate regions where the riblets exerted their most significant influence on cavitation. While the upstream-oriented SC-H riblets resulted in a smoother flow and delayed the transition to turbulence, the rotated SC-H riblets induced an unstable and unpredictable flow pattern. The images (a1), (a2), (a3), (a4) and (a5) show the cavitation dynamics at time steps $0 T$, $1/10 T$, $2/10 T$, $3/10 T$, $4/10 T$, respectively.

Figure 11 presents comparative plots of lift force amplitude (columns) versus cavitation shedding frequency (squares for 0 deg rotation, dots for 150 deg clockwise rotation) for the plain cylinder and for the cylinders equipped with horizontal and vertical riblets, caused by the flow at $\mathfrak{R} = 1.25 \times 10^5$ and $\sigma = 1.0$. Shown is the morphological variation for the 0 and 150 deg cylinder rotation, displaying lift force amplitude (columns) and vortex shedding frequency (squares and dots) for the plain cylinder, SC-H, SC-V, SA-H, and SA-V

riblets. Only the SC-H riblets caused a significant reduction of lift force amplitudes when rotated. In contrast, the other riblet configurations hardly affected the lift force amplitudes compared to those of the plain cylinder. Furthermore, the shedding frequency for the cylinders with the rotated riblets was similar to that for the plain cylinder.

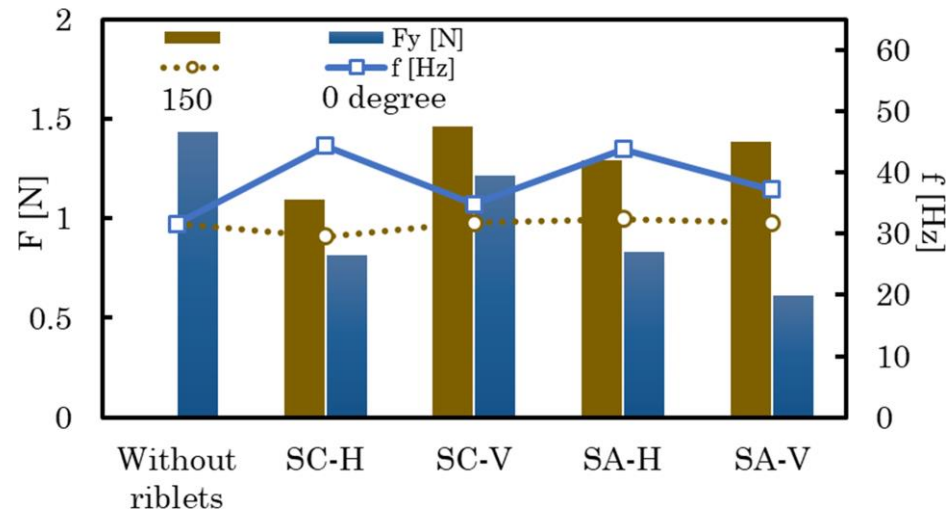


Figure 11. At $\Re = 1.25 \times 10^5$ and $\sigma = 1.0$, the morphological variations for the cylinder rotated 150 deg clockwise, displaying lift force amplitude (columns) and vortex shedding frequency (squares and dots) for the cylinder without riblets, with horizontal scalloped riblets (SC-H), with vertical scalloped riblets (SC-V), with horizontal sawtooth riblets (SA-H), and with vertical sawtooth riblets (SA-V). Only the SC-H riblets caused a significant reduction of lift force amplitudes when rotated. In contrast, the other riblet kinds hardly affected the lift force amplitudes compared to those for the smooth cylinder. Furthermore, rotating the riblets also hardly affected the shedding frequency, resembling the shedding frequency for the cylinder without riblets.

3.3. Effect of Reynolds Number

Cavitation effects may be more intense at higher Reynolds numbers, where the Tollmien–Schlichting (T–S) waves grow in magnitude and become more prominent. However, the impacts of riblets on cavitation are influenced by the specific flow conditions. Although they may still be effective in controlling cavitation at different Reynolds numbers, their effectiveness may be limited by stronger turbulence and pressure fluctuations associated with turbulent flows. Figure 12 shows cavitation structures behind the plain cylinder, SA-H, SA-V, SC-H, and SC-V riblets, which are caused by the flow at the lower $\Re = 1.0 \times 10^5$. As can be seen, the transition occurred further downstream, where flow separation from the cylinder's surface was less effective. Comparing Figure 12a with Figure 12b–e one can see that, at this lower $\Re = 1.0 \times 10^5$, both the SA-H and the SA-V riblets reduced cavitation by improving flow stability although the SA-V riblets were more effective in mitigating cavitation. The SC-H and the SC-V riblets also effectively inhibited cavitation although the SC-H riblets had a greater effect on mitigating cavitation. All riblets stabilized the flow, reduced pressure fluctuations, and led to more stable and smaller cavitation structures; however, their effect on cavitation patterns differed slightly.

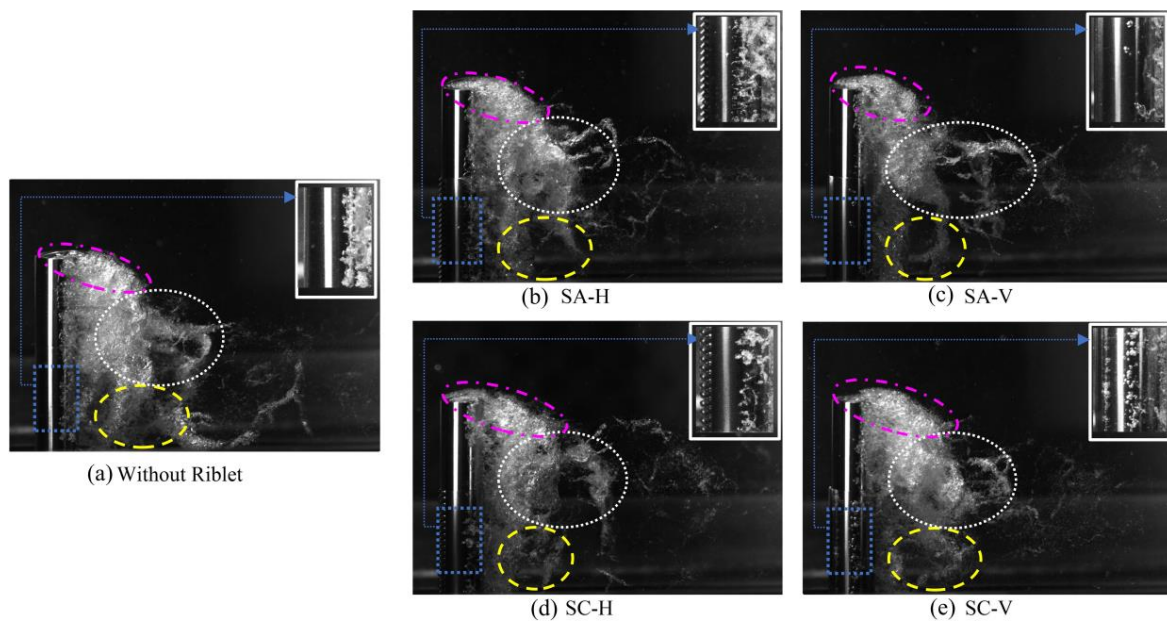


Figure 12. At the lower $\mathfrak{R} = 1.0 \times 10^5$ and $\sigma = 1.0$, cavity structures behind the plain cylinder (a), SA-H (b), SA-V (c), SC-H (d), and SC-V (e) cylinders.

Figure 13 shows five sequential images captured at regular intervals of the progression of cavitation around the cylinder with SC-H. This series of images, caused by the flow at the lower $\mathfrak{R} = 1.0 \times 10^5$, visualized the flow passing between riblet tips and valleys, thereby generating longitudinal vortices and facilitating effective vortex interaction within the boundary layer. Figure 14 shows cavitation dynamics around the plain cylinder, SA-H, with SA-V, SC-H, and SC-V riblets, which are caused by the flow at the higher $\mathfrak{R} = 1.5 \times 10^5$. As can be seen, now separation and flow reattachment occurred more readily due to the stronger turbulence generated by the T-S waves. Thus, the Reynolds number did influence the cavitation patterns. Comparative flow patterns shown in Figures 12a and 14a demonstrated that, for the cylinder without riblets caused by the flow at the lower Reynolds number, the cavitation patterns became less structured and less dense. While riblets were still effective in mitigating cavitation at higher Reynolds numbers, their effectiveness was diminished due to the increased flow energy.

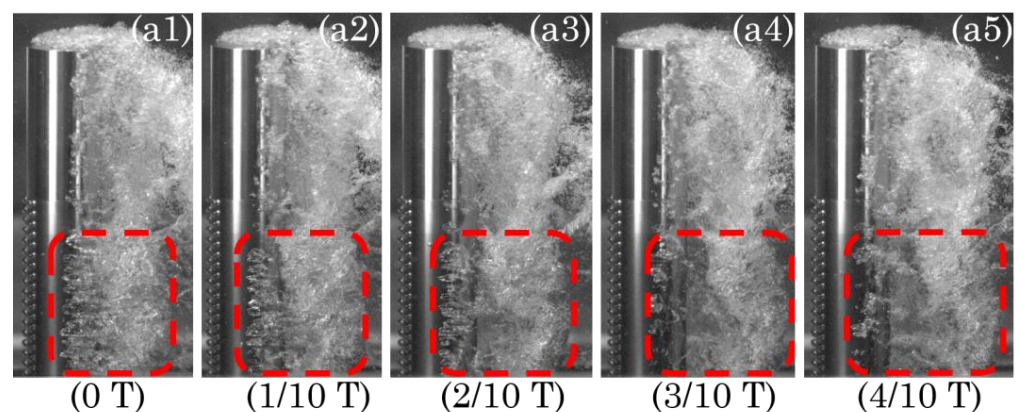


Figure 13. At the lower $\mathfrak{R} = 1.0 \times 10^5$ and $\sigma = 1.0$, five sequential images captured at regular intervals, showing the progression of cavitation behind a cylinder with SC-H riblets. These images were chosen from half of a standard cyclic period T that characterized cavitation shedding dynamics. Red surrounded areas indicate regions where the riblets exerted their most significant influence on cavitation. The images (a1), (a2), (a3), (a4) and (a5) show the cavitation dynamics at time steps $0 T$, $1/10 T$, $2/10 T$, $3/10 T$, $4/10 T$, respectively.

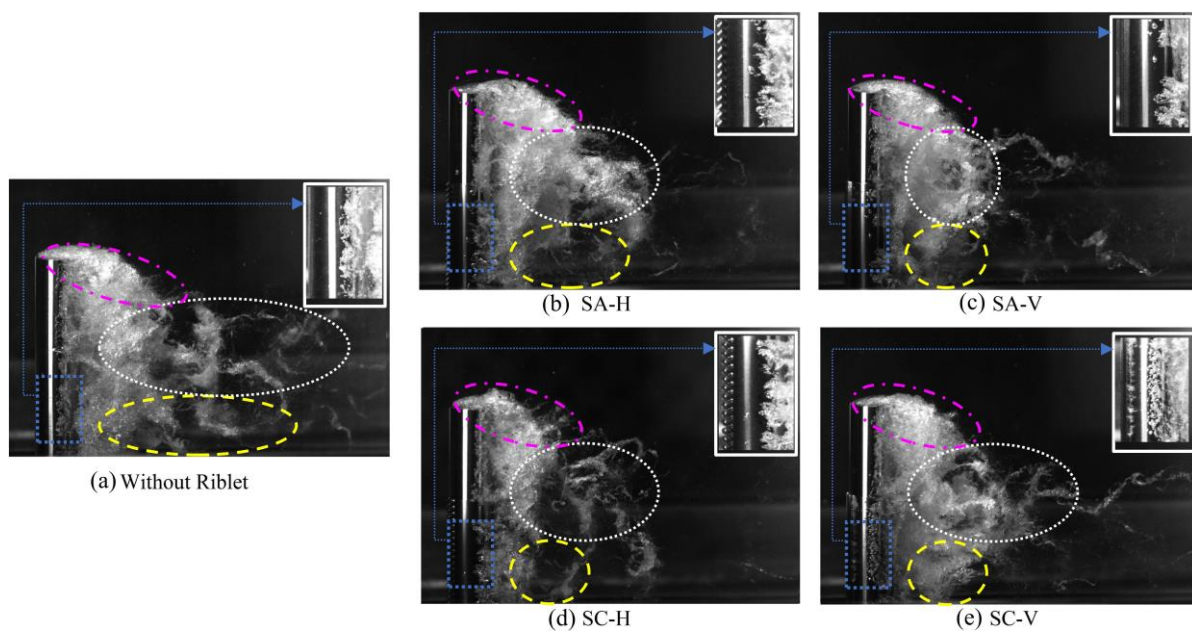


Figure 14. At the higher $\mathfrak{R} = 1.5 \times 10^5$ and $\sigma = 1.0$, cavity dynamics behind the plain cylinder (a), with SA-H (b), SA-V (c), SC-H (d), and SC-V (e) riblets.

Figure 14b–e show that, at this Reynolds number, SA-H, SA-V, and SC-V riblets provided only slight benefits although their effectiveness differed compared to the cases at lower Reynolds numbers. Notably, the SC-H riblets remained the most effective at mitigating the onset of cavitation. Furthermore, at $\mathfrak{R} = 1.5 \times 10^5$, cavitation occurred further upstream, and the cavitation patterns were more pronounced and elongated. Figure 15 presents five sequential images captured at regular intervals, showing the progression of cavitation dynamics behind a cylinder with SC-H riblets at higher Reynolds number $\mathfrak{R} = 1.5 \times 10^5$. Red surrounded areas indicate regions where the riblets exerted their most significant influence on cavitation. As can be seen, the SC-H riblets influenced the generation and distribution of the cavitation cloud, and they remained effective at delaying the onset of cavitation at this higher Reynolds number. Figure 16 plots comparative lift force amplitudes and cavitation shedding frequencies for cylinders with different kinds of riblets, caused by the flow at three different Reynolds numbers. Notably, at all Reynolds numbers, the SC-H riblets consistently reduced lift force amplitudes significantly. At $\mathfrak{R} = 1.0 \times 10^5$ and 1.25×10^5 , the other kinds of riblets also had a notable effect on lift force amplitudes. It was important to highlight that the introduction of riblets consistently led to an increase in shedding frequency at all Reynolds numbers. Interestingly, the shedding frequency of the cylinder with the SC-H riblets surpassed all other cases. This was due to the significantly reduced time required for cavities to detach themselves from the surface of the cylinder with SC-H riblets.

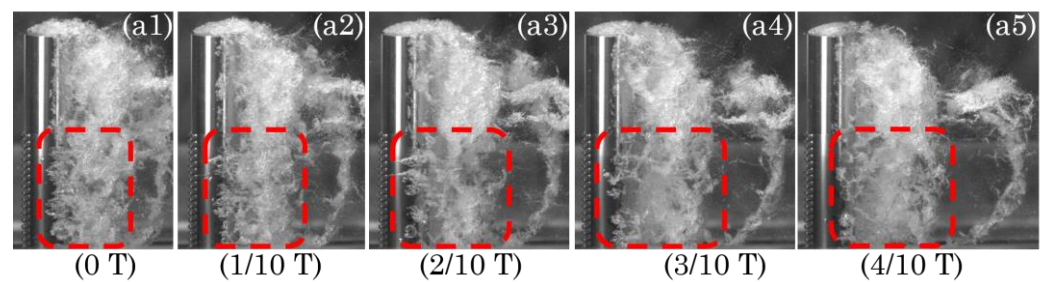


Figure 15. At the higher $\mathfrak{R} = 1.5 \times 10^5$ and $\sigma = 1.0$, five sequential images captured at regular intervals, showing the progression of cavitation behind a cylinder with SC-H riblets. These images were chosen from half of a standard cyclic period T that characterized cavity shedding dynamics. Red surrounded areas indicate regions where the riblets exerted their most significant influence on cavitation. The images (a1), (a2), (a3), (a4) and (a5) show the cavitation dynamics at time steps $0 T$, $1/10 T$, $2/10 T$, $3/10 T$, $4/10 T$, respectively.

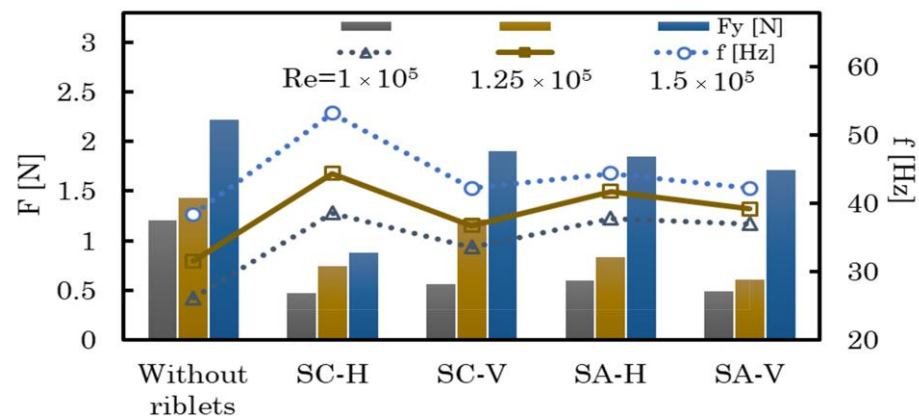


Figure 16. At $\sigma = 1.0$, lift force amplitude F and vortex shedding frequency f at $\mathfrak{R} = 1.0 \times 10^5$, 1.25×10^5 , and 1.5×10^5 for the plain cylinder and cylinders with SC-H, SC-V, SA-H, and SA-V riblets.

4. Conclusions

This study initially examined cavitation dynamics behind a smooth cylinder, obtaining high-speed images and measuring lift force fluctuations and shedding frequencies. Subsequently, the comprehensive investigation focused on cavitation control mechanisms using different mesoscale riblets inspired by biological structures; for example, sawtooth and scalloped riblets. The riblets, known for controlling flow separation and reducing drag, effectively generated streamwise vortices, influencing the formation of large-scale cavity structures. The key to their effectiveness lies in their ability to create vortices and control the boundary layer by transferring higher-momentum fluid to the low-momentum layer. Scalloped riblets induced secondary vortices near their tips, distorting primary streamwise vortices. Sawtooth riblets exhibited high velocity gradients near peaks but reduced gradients in valleys. Scalloped riblets, with their curved design, delayed turbulence transition, while sawtooth riblets mostly delayed flow separation. Horizontal riblets modified near-wall flow and turbulence, whereas vertical riblets created a smoother flow.

The considered riblets influenced flow dynamics and von Karman vortex characteristics, effectively reducing the cavitation-induced vibrations of the cylinder. SA-H and SC-H riblets notably decreased lift force fluctuations, reducing cavitation-induced vibration amplitudes by about 41% for SA-H and 43% for SC-H. Cavitation shedding frequencies increased to 43.8 Hz for SA-H and 44.4 Hz for SC-H compared to the smooth cylinder (31.6 Hz). Upstream SC-H riblets provided a smoother flow, delaying turbulence onset, while rotated SC-H riblets induced an unstable flow. The shedding frequency of rotated riblets was similar to the smooth cylinder. Riblets exhibit varying effectiveness in cavitation control at different Reynolds numbers. At lower Re (1.0×10^5), they enhance flow

stability and reduce cavitation, with SA-V riblets showing a pronounced effect. However, at higher Re (1.5×10^5), turbulence influences separation and reattachment, with SC-H riblets remaining the most effective at mitigating the onset of cavitation. The limitation of the passive control method such as riblets could be the high angle of attack of immersible bodies in cavitating regimes. At a high attack angle, a full flow separation may be achieved and, therefore, such a full separated flow cannot be affected using a passive control method. Therefore, this limitation should be considered in future works regarding implementation of the riblets on the hydrofoil surface and other industrial applications.

Author Contributions: Conceptualization, E.K.; methodology, E.K.; validation, E.K. and Y.L.; formal analysis, E.K., Y.L. and M.D.; investigation, E.K. and Y.L.; resources, E.K. and Y.L.; data curation, E.K., Y.L. and M.D.; writing—original draft preparation, E.K., Y.L. and M.D.; writing—review and editing, E.K., Y.L. and M.D.; visualization, E.K., Y.L. and M.D.; supervision, O.e.M.; project administration, O.e.M.; funding acquisition, E.K. and O.e.M. All authors have read and agreed to the published version of the manuscript.

Funding: This work was supported by the German Research Foundation with a Project Number of 469042952.

Data Availability Statement: The data supporting these findings were documented in this paper.

Acknowledgments: The design of the test cases was performed with the help of students Ataei and Ardalan.

Conflicts of Interest: The authors declare no conflicts of interest.

References

1. Reisman, G.; Wang, Y.; Brennen, C. Observations of shock waves in cloud cavitation. *J. Fluid Mech.* **1998**, *355*, 255–283. [\[CrossRef\]](#)
2. Dular, M.; Bachert, B.; Stoffel, B.; Sirok, B. Relationship between cavitation structures and cavitation damage. *Wear* **2004**, *257*, 1176–1184. [\[CrossRef\]](#)
3. Haosheng, C.; Yongjian, L.; Darong, C.; Jiadao, W. Experimental and numerical investigations on development of cavitation erosion pits on solid surface. *Tribol. Lett.* **2007**, *26*, 153159. [\[CrossRef\]](#)
4. Patella, R.; Choffat, T.; Reboud, J.; Archer, A. Mass loss simulation in cavitation erosion: Fatigue criterion approach. *Wear* **2013**, *300*, 205–215. [\[CrossRef\]](#)
5. Kadivar, E.; el Moctar, O.; Sagar, H. Experimental study of the influence of mesoscale surface structuring on single bubble dynamics. *Ocean Eng.* **2022**, *260*, 111892. [\[CrossRef\]](#)
6. Lin, Y.; Kadivar, E.; el Moctar, O.; Neugebauer, J.; Schellin, T.E. Experimental investigation on the effect of fluid–structure interaction on unsteady cavitating flows around flexible and stiff hydrofoils. *J. Phys. Fluids* **2022**, *34*, 083308. [\[CrossRef\]](#)
7. Franc, J.; Micheal, J. *Fundamentals of Cavitation*; Kluwer Academic Publishers: Dordrecht, The Netherlands, 2005.
8. Kadivar, E.; el Moctar, O.; Skoda, R.; Löschner, U. Experimental study of the control of cavitation-induced erosion created by collapse of single bubbles using a micro structured riblet. *Wear* **2021**, *486–487*, 204087. [\[CrossRef\]](#)
9. Arndt, R.E.; Ippen, A.T. *Cavitation near Surfaces of Distributed Roughness*; Massachusetts Institute of Technology: Cambridge, MA, USA, 1967.
10. Phan, T.-H.; Kadivar, E.; Nguyen, V.-T.; el Moctar, O.; Park, W.-G. Thermodynamic effects on single cavitation bubble dynamics under various ambient temperature conditions. *Phys. Fluids* **2022**, *34*, 023318. [\[CrossRef\]](#)
11. Venning, J.A.; Pearce, B.W.; Brandner, P.A. Nucleation effects on cloud cavitation about a hydrofoil. *J. Fluid Mech.* **2022**, *947*, A1. [\[CrossRef\]](#)
12. Yuan, G.Y.; Ni, B.Y.; Wu, Q.G.; Xue, Y.Z.; Han, D.F. Ice breaking by a high-speed water jet impact. *J. Fluid Mech.* **2022**, *934*, A1. [\[CrossRef\]](#)
13. Yuan, G.; Ni, B.; Wu, Q.; Lu, W.; Xue, Y. Experimental study on ice breaking by a cavitating water jet in a Venturi structure. *Appl. Therm. Eng.* **2024**, *239*, 122095. [\[CrossRef\]](#)
14. Kadivar, E. Experimental and Numerical Investigations of Cavitation Control Using Cavitating-Bubble Generators. Ph.D. Thesis, University of Duisburg-Essen, Duisburg, Germany, 2020.
15. Li, Q.; Franc, J.; Michel, J. Partial Cavities: Global Behavior and Mean Pressure. *J. Fluids Eng.* **1993**, *115*, 243–248.
16. Callenaere, M.; Franc, J.; Michel, J.; Rionde, M. The cavitation instability induced by the development of a re-entrant jet. *J. Fluid Mech.* **2001**, *444*, 223–256. [\[CrossRef\]](#)
17. Stutz, B.; Reboud, J.-L. Two-phase flow structure of sheet cavitation. *Phys. Fluids* **1997**, *9*, 3678–3686. [\[CrossRef\]](#)
18. Leroux, J.; Coutier-Delgosha, O.; Astolfi, A. A joint experimental and numerical study of mechanism associated to instability of partial cavitation on two-dimensional hydrofoil. *Phys. Fluids* **2005**, *17*, 052101. [\[CrossRef\]](#)
19. Leroux, J.; Astolfi, J.; Billard, J. An Experimental Study of Unsteady Partial Cavitation. *J. Fluids Eng.* **2004**, *126*, 94–101. [\[CrossRef\]](#)

20. Ganesh, H.; Mäkiharju, S.; Ceccio, S. Bubbly shock propagation as a mechanism for sheet-to-cloud transition of partial cavities. *J. Fluid Mech.* **2016**, *802*, 37–78. [\[CrossRef\]](#)
21. Ghahramani, E.; Ström, H.; Bensow, R.E. Numerical simulation and analysis of multi-scale cavitating flows. *J. Fluid Mech.* **2021**, *922*, A22. [\[CrossRef\]](#)
22. Kadivar, E.; Ochiai, T.; Iga, Y.; el Moctar, O. An experimental investigation of transient cavitation control on a hydrofoil using hemispherical vortex generators. *J. Hydrodyn.* **2020**, *33*, 1139–1147. [\[CrossRef\]](#)
23. Kadivar, E.; Timoshevskiy, M.; Nichik, M.; el Moctar, O.; Schellin, T.; Pervunin, K. Control of unsteady partial cavitation and cloud cavitation in marine engineering and hydraulic systems. *Phys. Fluids* **2020**, *32*, 052108. [\[CrossRef\]](#)
24. Long, X.; Cheng, H.; Ji, B.; Arndt, R.E.A.; Peng, X. Large eddy simulation and Euler–Lagrangian coupling investigation of the transient cavitating turbulent flow around a twisted hydrofoil. *Int. J. Multiph. Flow* **2018**, *100*, 41–56. [\[CrossRef\]](#)
25. Pelz, P.; Keil, T.; Groß, T.F. The transition from sheet to cloud cavitation. *J. Fluid Mech.* **2017**, *817*, 439–454. [\[CrossRef\]](#)
26. Wang, Z.; Cheng, H.; Ji, B. Euler–Lagrange study of cavitating turbulent flow around a hydrofoil. *Phys. Fluids* **2021**, *33*, 112108. [\[CrossRef\]](#)
27. Wang, Z.; Cheng, H.; Ji, B.; Peng, X. Numerical investigation of inner structure and its formation mechanism of cloud cavitating flow. *Int. J. Multiphase Flow* **2023**, *165*, 104484. [\[CrossRef\]](#)
28. Matsudaira, Y.; Gomi, Y.; Oba, R. Characteristics of bubble-collapse pressures in a karman-vortex cavity. *JSME Int. J.* **1992**, *35*, 179–185. [\[CrossRef\]](#)
29. Saito, Y.; Sato, K. Cavitation bubble collapse and impact in the wake of a circular cylinder. In Proceedings of the Fifth International Symposium on Cavitation, Osaka, Japan, 1–4 November 2003; pp. 3–8.
30. Kumar, P.; Chatterjee, D.; Bakshi, S. Experimental investigation of cavitating structures in the near wake of a cylinder. *Int. J. Multiph. Flow* **2017**, *89*, 207–217. [\[CrossRef\]](#)
31. Geike, T. Bubble dynamics-based modeling of the cavitation dynamics in lubricated contacts. *Facta Univ. Ser. Mech. Eng.* **2021**, *19*, 115–124. [\[CrossRef\]](#)
32. Gnanaskandan, A.; Mahesh, K. Numerical investigation of near-wake characteristics of cavitating flow over a circular cylinder. *J. Fluid Mech.* **2016**, *790*, 453–491. [\[CrossRef\]](#)
33. Dobroselsky, K. Cavitation streamlining of a round cylinder in the critical range. *J. Phys. Conf. Ser.* **2020**, *1677*, 012056. [\[CrossRef\]](#)
34. Brandao, F.; Bhatt, M.; Mahesh, K. Numerical study of cavitation regimes in flow over a circular cylinder. *J. Fluid Mech.* **2019**, *885*, A19. [\[CrossRef\]](#)
35. Gu, F.; Huang, Y.; Zhang, D. Cavitation of Multiscale Vortices in Circular Cylinder Wake at $Re = 9500$. *J. Mar. Sci. Eng.* **2021**, *9*, 1366. [\[CrossRef\]](#)
36. Sadri, M.; Kadivar, E. Numerical investigation of the cavitating flow and the cavitation-induced noise around one and two circular cylinders. *Ocean. Eng.* **2023**, *277*, 114178. [\[CrossRef\]](#)
37. Zhang, A.M.; Li, S.-M.; Cui, P.; Li, S.; Liu, Y.-L. A unified theory for bubble dynamics. *Phys. Fluids* **2023**, *35*, 033323. [\[CrossRef\]](#)
38. Custodio, D.; Henoch, C.; Johari, H. Cavitation on hydrofoils with leading edge protuberances. *Ocean Eng.* **2018**, *162*, 196–208. [\[CrossRef\]](#)
39. Hao, J.; Zhang, M.; Huang, X. The influence of surface roughness on cloud cavitation flow around hydrofoils. *Acta Mech. Sin.* **2018**, *34*, 10–21. [\[CrossRef\]](#)
40. Zhao, W.; Wang, G. Research on passive control of cloud cavitation based on a bionic fin-fin structure. *Eng. Comput.* **2020**, *37*, 863–880. [\[CrossRef\]](#)
41. Zhang, L.; Chen, M.; Shao, X. Inhibition of cloud cavitation on a flat hydrofoil through the placement of an obstacle. *Ocean Eng.* **2018**, *155*, 1–9. [\[CrossRef\]](#)
42. Che, B.; Chu, N.; Schmidt, S.; Cao, L.; Likhachev, D.; Wu, D. Control effect of micro vortex generators on leading edge of attached cavitation. *Phys. Fluids* **2019**, *31*, 044102-1–044102-12. [\[CrossRef\]](#)
43. Che, B.; Chu, N.; Cao, L.; Schmidt, S.; Likhachev, D.; Wu, D. Control effect of micro vortex generators on attached cavitation instability. *Phys. Fluids* **2019**, *31*, 064102-1–064102-11. [\[CrossRef\]](#)
44. Kadivar, E.; Timoshevskiy, M.; Pervunin, K.; el Moctar, O. Cavitation control using cylindrical cavitating-bubble generators (CCGS): Experiments on a benchmark cav2003 hydrofoil. *Int. J. Multiph. Flow* **2020**, *125*, 103186. [\[CrossRef\]](#)
45. Kadivar, E.; el Moctar, O. Investigation of cloud cavitation passive control method for hydrofoils using Cavitating-bubble Generators (CGs). In Proceedings of the 10th International Symposium on Cavitation, Baltimore, MD, USA, 14–16 May 2018.
46. Simanto, R.; Hong, J.; Kim, K.; Ahn, B.; Shin, S. Experimental investigation on cavitation and induced noise of two-dimensional hydrofoils with leading-edge protuberances. *Phys. Fluids* **2022**, *34*, 124115. [\[CrossRef\]](#)
47. Yu, F.; Zhang, Y.; Liu, H.; Zhou, Q. Experimental and numerical investigation of cavitation control with porous material on the hemisphere cylinder. *J. Ocean. Eng.* **2022**, *266*, 112984. [\[CrossRef\]](#)
48. Kawamura, T.; Hiwada, M.; Hibino, T.; Mabuchi, I.; Kumada, M. Flow around a Finite Circular Cylinder on a Flat Plate: Cylinder height greater than turbulent boundary layer thickness. *Bull. JSME* **1984**, *27*, 2142–2151. [\[CrossRef\]](#)
49. Kitagawa, T.; Fujino, Y.; Kimura, K.; Mizuno, Y. Wind pressures measurement on end-cell-induced vibration of a cantilevered circular cylinder. *J. Wind Eng. Ind. Aerodyn.* **2002**, *90*, 395–405. [\[CrossRef\]](#)
50. Graf, W.H.; Yulistiyanto, B. Experiments on Flow around a cylinder; the velocity and vorticity fields. *J. Hydraul. Res.* **1998**, *36*, 637–653. [\[CrossRef\]](#)

51. Park, C.W.; Lee, S.J. Free end effects on the near wake flow structure behind a finite circular cylinder. *J. Wind Eng. Ind. Aerodyn.* **2000**, *88*, 231–246. [[CrossRef](#)]
52. van Rijsbergen, M. A review of sheet cavitation inception mechanisms. In Proceedings of the IS-ROMAC2016—International Symposium on Transport Phenomena and Dynamics of Rotating Machinery, Honolulu, HI, USA, 10–16 April 2016.
53. Bixler, G.D.; Bhushan, B. Fluid drag reduction with shark-skin riblet inspired microstructured surfaces. *J. Adv. Funct. Mater.* **2013**, *23*, 4507–4528. [[CrossRef](#)]
54. Lulekar, S.S.; Ghassemi, P.; Alsalihi, H.; Chowdhury, S. Adaptive-fidelity design automation framework to explore bioinspired surface riblets for drag reduction. *AIAA J.* **2021**, *59*, 880–892. [[CrossRef](#)]
55. Wang, Y.; Huang, Y.; Fu, S. On the tip sharpness of riblets for turbulent drag reduction. *J. Acta Mech. Sin.* **2022**, *38*, 321389. [[CrossRef](#)]
56. Lee, S.J.; Lee, S.H. Flow field analysis of a turbulent boundary layer over a riblet surface. *Exp. Fluids* **2001**, *30*, 153–166. [[CrossRef](#)]
57. Goldstein, D.; Handler, R.; Sirovich, L. Direct numerical simulation of turbulent flow over a modelled riblet covered surface. *J. Fluid Mech.* **1995**, *302*, 333–376. [[CrossRef](#)]
58. Dean, B.; Bharat, B. Shark-skin surfaces for fluid-drag reduction in turbulent flow: A review. *Phil. Trans. R. Soc. A* **2010**, *368*, 4775–4806. [[CrossRef](#)] [[PubMed](#)]

Disclaimer/Publisher’s Note: The statements, opinions and data contained in all publications are solely those of the individual author(s) and contributor(s) and not of MDPI and/or the editor(s). MDPI and/or the editor(s) disclaim responsibility for any injury to people or property resulting from any ideas, methods, instructions or products referred to in the content.